Evolution of viscous conducting drops subject to rotation and electric fields

## V. J. García-Garrido

#### joint work with M. A. Fontelos and U. Kindelán

12<sup>th</sup> International Conference on Free Boundary Problems Theory and Applications Germany 11 - 15 June 2012





 Problem
 Numerical Method
 Results
 Conclusions & Future Research Lines

 •0
 00
 000000000
 000000000

# Mathematical Model



#### **Two cases:** $L = \mathcal{I}\omega$

- Electric field  $\mathcal{E}_{\infty}$  and constant L.
- Charge Q and constant L.

#### Notation

- $\begin{aligned} \mathcal{D}_1(t) \ \text{Fluid drop domain} \\ \rho_1 \ \text{Drop density} \\ \mu_1 \ \text{Drop viscosity} \\ p^{(1)} \ \text{Drop pressure} \\ \mathbf{u}^{(1)} \ \text{Drop velocity} \end{aligned}$
- $\mathcal{D}_2(t)$  Outer fluid domain  $\rho_2$  Outer fluid density  $\mu_2$  Outer fluid viscosity  $p^{(2)}$  Outer fluid pressure  $\mathbf{u}^{(2)}$  Outer fluid velocity

 $\boldsymbol{\omega}$  Angular velocity  $\partial \mathcal{D}(t)$  Fluid interface

Problem ⊙●	Numerical Method	Results 00000000	Conclusions & Future Research Line
Mathematic	cal Model		

## Stokes equation

$$\begin{cases} \mu_i \Delta \mathbf{u}^{(i)} - \nabla \Pi^{(i)} = \mathbf{0} &, & \text{in } \mathcal{D}_i(t) \\ \nabla \cdot \mathbf{u}^{(i)} = 0 &, & \text{in } \mathcal{D}_i(t) \end{cases}, \quad i \in \{1, 2\}$$

#### Boundary condition

$$\left(T^{(2)} - T^{(1)}\right)\mathbf{n} = \left(2\gamma \mathcal{H} - \varrho \frac{L^2}{2\mathcal{I}^2}r_{axis}^2 - \frac{\sigma^2}{2\varepsilon_0}\right)\mathbf{n} \quad , \quad \text{on } \partial \mathcal{D}\left(t\right)$$

## Laplace equation

$$\begin{cases} \Delta \mathcal{V} = 0 & , \quad \text{in } \mathcal{D}_2\left(t\right) \\ \mathcal{V} = \mathcal{V}_0 & , \quad \text{in } \partial \mathcal{D}\left(t\right) \\ \mathcal{V} \to -\mathcal{E}_{\infty} z + O(|\mathbf{r}|^{-1}) & , \quad \text{as } |\mathbf{r}| \to \infty \end{cases}$$

 Problem
 Numerical Method
 Results
 Conclusions & Future Research Lines

 00
 ●0
 000000000

# Boundary integral formulation

Fredholm integral equation of the  $2^{nd}$  kind for the velocity:

$$\begin{split} u_{j}(\mathbf{r}') &= -\frac{1}{4\pi \left(\mu_{1} + \mu_{2}\right)} \underbrace{\int_{\partial \mathcal{D}(t)} f_{i}(\mathbf{r}) G_{ij}(\mathbf{r}, \mathbf{r}') \, dS(\mathbf{r})}_{\text{Single layer potential}} \\ &- \frac{\mu_{1} - \mu_{2}}{4\pi \left(\mu_{1} + \mu_{2}\right)} \underbrace{\int_{\partial \mathcal{D}(t)}^{PV} u_{i}(\mathbf{r}) T_{ijk}(\mathbf{r}, \mathbf{r}') n_{k}(\mathbf{r}) \, dS(\mathbf{r})}_{\text{Druble layer potential}} \end{split}$$

Double layer potential

$$\begin{split} G_{ij}(\mathbf{r}, \mathbf{r}') &= \frac{\delta_{ij}}{|\mathbf{r} - \mathbf{r}'|} + \frac{(r_i - r'_i)(r_j - r'_j)}{|\mathbf{r} - \mathbf{r}'|^3} & \text{Stokeslet} \\ T_{ijk}(\mathbf{r}, \mathbf{r}') &= -6 \frac{(r_i - r'_i)(r_j - r'_j)(r_k - r'_k)}{|\mathbf{r} - \mathbf{r}'|^5} & \text{Tenselet} \\ f_i(\mathbf{r}) &= \left[ 2\gamma \mathcal{H}(\mathbf{r}) - \varrho \frac{L^2}{2\mathcal{I}^2} r_{axis}^2 - \frac{\sigma^2}{2\varepsilon_0} \right] n_i(\mathbf{r}) & \text{Traction} \end{split}$$

where  $i,j,k\in\left\{ 1,2,3\right\}$  and  $\mathbf{r}$  ,  $\mathbf{r}^{\prime}\in\partial\mathcal{D}\left(t\right).$ 

Problem	Numerical Method	Results	Conclusions & Future Research Lines
00	○●	000000000	
Numerical a	lgorithm		

#### Algorithm

- **O** Compute the volume and moment of inertia about the *z*-axis.
- ② Calculate mean curvature and charge density of drop.
- Solve linear system to obtain velocity field at the boundary.
- Move the boundary with an Euler explicit scheme:

$$\mathbf{r}(t_{n+1}) = \mathbf{r}(t_n) + \mathbf{u}(t_n) \Delta t .$$

- Segularization of the mesh (if necessary):
  - Delaunay remeshing.
  - Mesh relaxation.
  - Mesh refinement.

Repeat above steps until  $t_{max}$  is reached.

Problem Numerical Method Results Conclusions & Future Research Lines

# Charge and rotation at constant L

## Asymptotic expansion: $L^2 \ll 1$

## Young-Laplace equation

$$\delta p = 2\gamma \mathcal{H} - \varrho \frac{L^2}{2\mathcal{I}^2} r_{axis}^2 - \frac{\sigma^2}{2\varepsilon_0} \quad , \quad \text{on } \partial \mathcal{D}$$

= 1

## Spheroidal approximation:

Energy formulation

Minimize: 
$$E_{total} = E_{area} + E_{kinetic} + E_{electrostatic}$$
, V



## Charge and rotation at constant L



 Problem
 Numerical Method
 Results
 Conclusions & Future Research Lines

 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00
 00

## Charge and rotation at constant L



 $\chi=0.892$ 

Problem Numerical Method Results Conclusions & Future Research Lines

## Charge and rotation at constant L

#### Stability analysis



Problem 00	Numerical Method	Results	Conclusions & Future Research Lines

## Electric field and rotation at constant L



 $E_{\infty}^{2} = h\left(\alpha\right) + g\left(\alpha\right)L^{2}$ 



## Electric field and rotation at constant L



Stability curve for axisymmetric shapes





Centrifugal force vs. Surface tension  $(L = 0.2, E_{\infty} = 0.9)$ 



#### Self-similarity





# Conclusions & Future Research Lines

#### Conclusions

- Adaptive BEM to simulate droplet evolution.
- Theoretical models to approximate charged rotating drops.
- Stability analysis shows ellipsoidal configurations and singularities (Taylor cones and two-lobed drop breakup).
- Linear relationship for shapes of rotating drops subject to uniform electric fields with same aspect ratio.
- Taylor cone semiangle is not affected by small rotations.

#### Research Lines

- Evolution of charged rotating drops subject to electric fields.
- Describe the space of parameters  $(\chi, E_{\infty}, L)$ .
- Understand the role of rotation on the stability of the system.

Problem	Numerical Method	Results	Conclusio

# Thank you for your attention.

# Questions?

Problem	Numerical Method	Results	Conclusions & Future Research Lines
00	00	000000000	
Mean curva	ture		

## Algorithm 1 Paraboloid fitting

- 1: Take  $\mathbf{n}_p$  initial approximation to the outward unit normal to  $\partial \mathcal{D}$  at  $\mathbf{p}$ .
- 2: repeat
- 3: Choose local coordinates (x', y', z') at  $\mathbf{p}$  and z'-axis along  $\mathbf{n}_p$ .
- 4: Find  $(x'_i, y'_i, z'_i)$  coordinates of the adjacent nodes to **p**.
- 5: Minimize:

$$F = \sum_{i=1}^{N_p} \left( Ax'_i + By'_i + C(x'_i)^2 + Dx'_i y'_i + E(y'_i)^2 - z'_i \right)^2 \,.$$

$$\begin{split} & \textbf{6:} \qquad (\mathbf{n}_p)_n \leftarrow (-A, -B, 1)/(1+A+B)^{1/2}. \\ & \textbf{7:} \quad \textbf{until} \ |(\mathbf{n}_p)_n - \mathbf{n}_p| < \varepsilon. \\ & \textbf{8:} \ \text{Mean curvature} \ k_{\mathbf{p}} = \frac{(1+B^2)C - ABD + (1+A^2)E}{(1+A^2+B^2)^{3/2}} \end{split}$$

#### Fast convergence iterative method

Problem	Numerical Method	Results	Conclusions & Future Research Lines
00	00	00000000	
C C I	1		

# Surface charge density

Write:

$$\sigma = \mathcal{V}_0 \sigma_0 + \sigma_{ind}$$

Solve:

$$\mathcal{E}_{\infty} z_{0} = \frac{1}{4\pi\varepsilon_{0}} \int_{\partial \mathcal{D}(t)} \frac{\sigma_{ind}\left(\mathbf{x}\right)}{\|\mathbf{x} - \mathbf{x}_{0}\|} \, dS \quad , \quad 1 = \frac{1}{4\pi\varepsilon_{0}} \int_{\partial \mathcal{D}(t)} \frac{\sigma_{0}\left(\mathbf{x}\right)}{\|\mathbf{x} - \mathbf{x}_{0}\|} \, dS$$

with:

$$Q = \mathcal{V}_0 \int_{\partial \mathcal{D}(t)} \sigma_0 \left( \mathbf{x} \right) \ dS + \int_{\partial \mathcal{D}(t)} \sigma_{ind} \left( \mathbf{x} \right) \ dS$$

Linear system at barycenters  $\mathbf{x}_l$ ,  $\mathbf{x}_j$  of triangles:

$$\int_{\partial \mathcal{D}(t)} \frac{\sigma(\mathbf{x})}{\|\mathbf{x} - \mathbf{x}_i\|} \, dS \approx \sum_{j=1}^M \lambda_{ij} \sigma_j \quad , \quad \lambda_{ij} = \int_{T_j} \frac{dS}{\|\mathbf{x} - \mathbf{x}_i\|}$$
$$\sigma_j = \sigma(\mathbf{x}_j)$$

Problem 00	Numerical Method	Results 00000000	Conclusions & Future Research Lines	
Surface charge density				







• Case 
$$i = j$$

$$\lambda_{ii} = \sum_{k=1}^{6} \int \int_{T_{ik}} d\rho \, d\theta = \sum_{k=1}^{6} a_k \ln\left(\sec\left(\alpha_k\right) + \tan\left(\alpha_k\right)\right)$$

• Case  $i \neq j$ 

$$\lambda ij = \sum_{k=1}^{N_s} \lambda_{ij,k} \quad , \quad \lambda_{ij,k} = \frac{Area\left(T_{jk}\right)}{\|\mathbf{b}_{jk} - \mathbf{x}_i\|}$$