



Weierstrass Institute for  
Applied Analysis and Stochastics



Leibniz Institute for  
Crystal Growth



# Modeling and numerical simulation of the application of traveling magnetic fields to stabilize crystal growth from the melt

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joined work with:

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Ch. Frank-Rotsch, F.-M. Kießling, W. Miller, U. Rehse, (IKZ, Berlin),

P. Rudolph (CTC, Schönefeld)



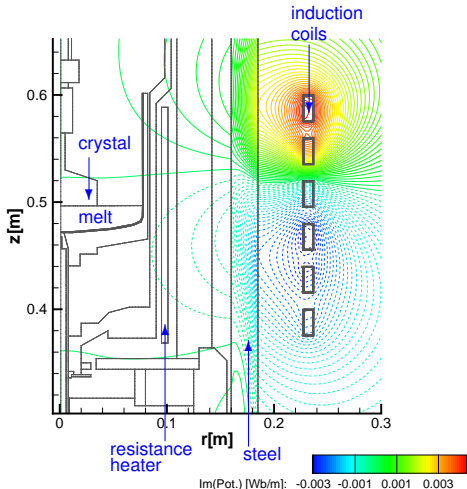
DFG Research Center MATHEON  
Mathematics for key technologies

## Magnetic fields and crystal growth processes

- Semi-conductor mono-crystals are important for high-technology devices
- Many of the crystals: grown from the melt by growth processes of Czochralski type
- To improve the growth processes, one needs to stabilize the melt movement
- (Time-dependent) magnetic field can stabilize the melt movement
- Typically, the magnetic fields are generated by induction coils placed outside of the growth apparatus
- Growth of III-V compounds requires to use pressure chambers with thick steel walls  $\rightsquigarrow$  walls diminish the magnetic field generated by external coils

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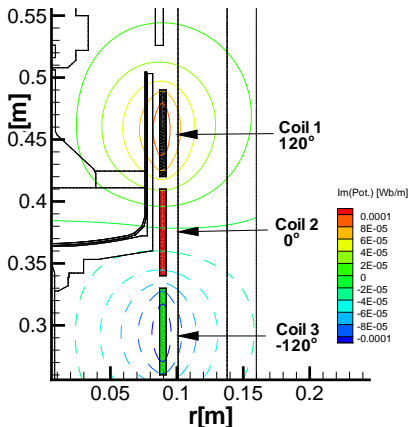
$\rightsquigarrow$  producing a field of sufficient magnitude in the melt requires much energy

- **Cooperation of :**

- Leibniz Institute of Crystal Growth (IKZ)
  - ETP (Hannover)
  - IISB (Erlangen)
- WIAS
- Steremat Elektrowärme GmbH, Berlin
- AUTEAM Industrie-Elektronik GmbH, Brandenburg

- Internal *heater-magnet modules (HMM)*, i.e. coil-formed resistance heaters, and electrical components to use them have been developed

- Replacing the usual meander-formed resistance heater units in the growth vessel by an HMM
  - ↪ one can generate appropriate fields in the melt with moderate power consumption



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**Innovation award Berlin-Brandenburg 2008**  
given to the project

- Navier-Stokes equations in Boussinesq approximation: for melt in  $\Omega_1$

$$\rho_1 \left( \frac{\partial u}{\partial t} + (u \cdot \nabla) u \right) = -\nabla p + \operatorname{div}(2\eta(\theta) D u) + f(\theta) + j \times \mu H,$$

$$\operatorname{div} u = 0 \quad \text{in } ]0, T[ \times \Omega_1 .$$

- Maxwell's equations

$$j = \operatorname{curl} H = 0, \quad \operatorname{div}(\sigma_c E) = 0, \quad \text{in } ]0, T[ \times \tilde{\Omega}_{nc},$$

$$j = \operatorname{curl} H = \sigma_c(\theta) (E + u \times \mu H) \quad \text{in } ]0, T[ \times \tilde{\Omega}_c$$

$$\operatorname{curl} E + \mu \frac{\partial H}{\partial t} = 0, \quad \operatorname{div}(\mu H) = 0, \quad \text{in } ]0, T[ \times \tilde{\Omega} .$$

- Energy balance

$$\rho_1 c \left( \frac{\partial \theta}{\partial t} + u \cdot \nabla \theta \right) = \operatorname{div}(\kappa(\theta) \nabla \theta) + \eta(\theta) D(u, u) + \frac{|j|^2}{\sigma_c(\theta)} \quad \text{in } ]0, T[ \times \Omega_1,$$

$$\rho c \frac{\partial \theta}{\partial t} = \operatorname{div}(\kappa(\theta) \nabla \theta) + \frac{|j|^2}{\sigma_c(\theta)} \quad \text{in } ]0, T[ \times \Omega_i \quad (i \neq 1) .$$



$$-\kappa_{\text{gas}}(\theta_{\text{gas}})\nabla\theta_{\text{gas}} \bullet \vec{n}_{\text{gas}} - R + J = -\kappa(\theta)\nabla\theta \bullet \vec{n}_{\text{gas}}$$

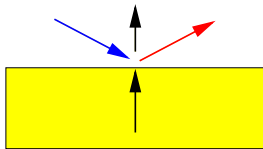
- $\vec{n}_{\text{gas}}$ : outer unit normal w.r.t. gas phase,
- total outgoing radiation  $R = \sigma\epsilon T_{\text{gl0}}^4 + (1 - \epsilon)J$ 
  - $\sigma$  Boltzmann radiation constant,  $\epsilon$  emissivity

- incoming radiation

$$J(x) = \int_{\Gamma} \Lambda(x, y) \omega(x, y) R(y) dy$$

- $\Gamma$  boundary of cavity
- $\Lambda(x, y) = 1$  if  $y$  is "visible" from  $x$ , 0 otherwise
- $\omega(x, y)$  view factors

$$\omega(x, y) := \frac{(\vec{n}_{\text{gas}}(y) \bullet (x - y)) (\vec{n}_{\text{gas}}(x) \bullet (y - x))}{\pi((y - x) \bullet (y - x))^2}$$



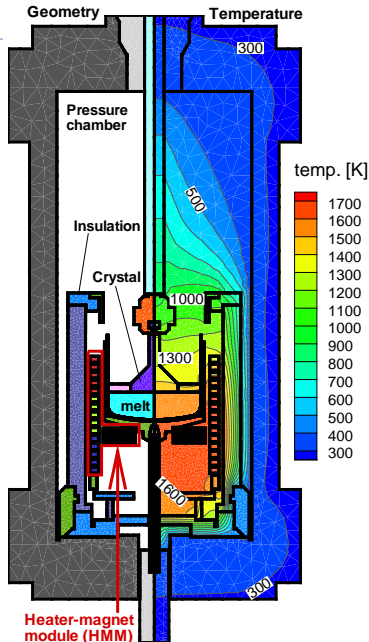
- Concept of HMM: Rudolph JCG 2008
- Modeling and Simulation: Lechner–K.–Druet JCG 2007  
K.–Lechner–Druet–Philip–Sprekels–Frank–Rotsch–Kießling–Miller–Rehse–Rudolph JCG 2008, MHD 2009  
Rudolph–Czupalla–Dropka–Frank–Rotsch–Kiessling–K.-Lux–Miller–Rehse–Root JKCGC 2009  
Dropka–Miller–Rehse–Rudolph–Buellesfeld–Sahr–K. –Reinhardt JCG 2011  
Dreyer–Druet–K.–Sprekels WIAS 2012
- Existence of solutions: Druet Thesis 2009, MMAS 2009, CzMJ 2009, NA-RWA 2009, ApM 2010,
- Optimal control problem: Druet–K.–Sprekels–Tröltzsch–Yousept SIAM JCO 2011
- Free Boundary Problem: Druet WIAS 2011, WIAS 2012,

**Contributed Talk Wed. 11.00**

P.-É. Druet is researcher in the project C9 “Simulation and Optimization of Semiconductor Crystal Growth from the Melt Controlled by Traveling Magnetic Fields” in the DFG-Research Center FZT 86, MATHEON, (Heads: O. Kl., J. Sprekels, F. Tröltzsch)

## Global simulations

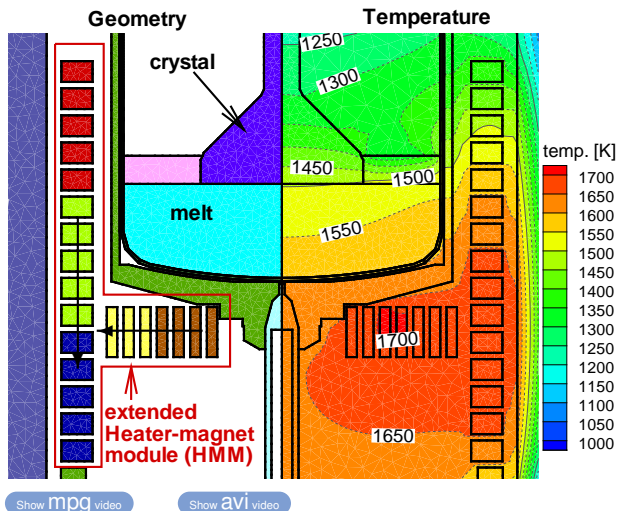
- main parts of crystal growth device are simulated, axially symmetric approximation of real geometry is considered
- melt motion is ignored
- time-harmonic version Maxwell equation
- stationary heat equation
- Software **WIAS-HiTNIHS** (P. Philip, O. Kl.)  
**WIAS-High Temperature Numerical Induction Heating Simulator** (partially developed in MATHEON Project C9)
- LPA Mark 3 in a configuration for LEC crystal growth of GaAs  
4 kg GaAs melt, diameter=15.2 cm, height=4.5 cm  
GaAs melt covered by a boric oxide layer with a height of 1.35 cm  
Control T at triple point by adapting power used in simulation



Special HMM following the patent DE 10 2007 028 548 by Ch. Frank-Rotsch, P. Rudolph, O. Kl., R.-P.

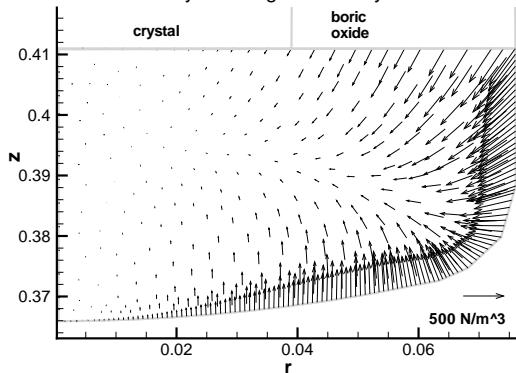
Lange, B. Nacke:

- 3 coils surrounding the crucible, each having 5 rings, producing a downwards moving TMF,
- 2 additional spirals below the crucible, producing an outward moving TMF





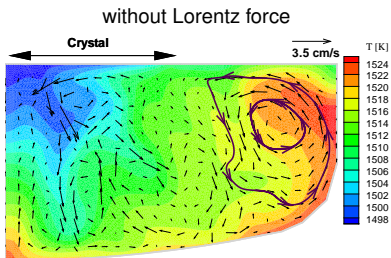
- time average over one period of electromagnetic fields, computed from time-harmonic representation  $\mapsto$  **Lorentz-force density for local simulation**
- Lorentz force density in melt generated by the HMM in the LPA Mark 3:



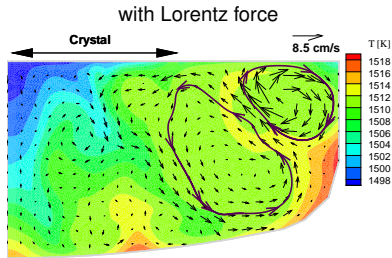
- Temperature and heat fluxes in the melt computed by **global** simulations are also used as **initial and boundary data for local simulation in the melt**

- only domain filled with melt is considered.
- heat sources by induction current ignored
- motion induced current  $\vec{u} \times \vec{B} = \vec{u} \times \mu \vec{H}$  neglected  $\rightsquigarrow$  Maxwell equation are decoupled from melt motion
- Implementation by Ch. Lechner in the framework of NAVIER (E. Bänsch)

Snapshots of computed velocity and temperature distribution

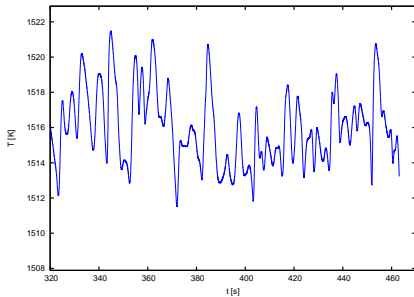


maximal velocity  $\approx 3.5 \frac{\text{cm}}{\text{s}}$



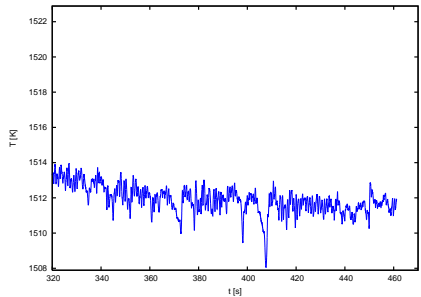
maximal velocity  $\approx 8.5 \frac{\text{cm}}{\text{s}}$

## Temperature oscillations in a monitor point (Ch. Lechner using NAVIER)



without Lorentz force

main temperature oscillations  $\approx 8\text{K}$



with Lorentz force

main temperature oscillations  $\approx 2\text{K}$   
frequency increased

- Conclusion project KRISTMAG<sup>®</sup> :
  - Numerical simulations + crystal growth experiments show
    - Lorentz forces generated by an internal HMM can influence the melt flow during crystal growth
    - using an extended HMM and appropriate TMFs, we can improve the growth conditions
  - an extended HMM has successfully been used for LEC crystal growth at the IKZ
- Project AVANTSOLAR (2008-2011):
  - KRISTMAG<sup>®</sup> project partners + SCHOTT Solar Wafer GmbH + two other research institutes
  - using traveling magnetic fields generated by an internal HMM to improve directional solidification of solar-grad silicon
  - successfully growth of 640 kg Si ingots

- Projects **KRISTMAG**<sup>®</sup> (07/2005–06/2008) and **AVANT SOLAR** (07/2008–06/2011) were supported by
  - the German Federal State of Berlin in the framework of the “Zukunftsfonds Berlin”,
  - the Technology Foundation Innovation Center Berlin (TSB),
  - cofinanced by the European Union within the European Regional Development Fund (EFRE). Investing in your future.



- Project C9 “Simulation and Optimization of Semiconductor Crystal Growth from the Melt Controlled by Traveling Magnetic Fields” (05/2002–05/2014) is supported by:



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