

Finite Element Techniques for the Simulation of Two-phase Incompressible Flows

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Free

Introduction:

- Standard **model** for two-phase flow.

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Nitsche's method for a mass transport problem:

- Explanation of the method.
- Results of numerical experiments.

Sharp Interface Model for two-phase flows with interfacial phenomena

Fluid dynamics in the two phases

Incompressible **Navier-Stokes** in the subdomains:

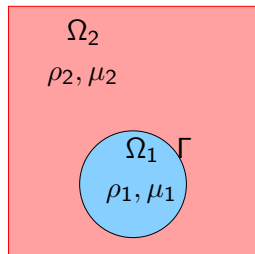
Domains: $\Omega_1 = \Omega_1(t)$, $\Omega_2 = \Omega_2(t)$

Interface: $\Gamma = \Gamma(t)$

ρ_i : density in Ω_i

μ_i : viscosity in Ω_i

$$\mathbf{D}(\mathbf{u}) = \nabla \mathbf{u} + \nabla \mathbf{u}^T, \quad \boldsymbol{\sigma} = -p \mathbf{I} + \mu \mathbf{D}(\mathbf{u})$$



$$\begin{cases} \rho_i(\mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u}) = \operatorname{div}(\boldsymbol{\sigma}) + \rho_i \mathbf{g} \\ \quad = -\nabla p + \operatorname{div}(\mu_i \mathbf{D}(\mathbf{u})) + \rho_i \mathbf{g} & \text{in } \Omega_i \quad \text{for } i = 1, 2 \\ \operatorname{div} \mathbf{u} = 0 & \text{in } \Omega_i \end{cases}$$

Interfacial effects

$$\text{Force on } \gamma = \int_{\partial\gamma} \boldsymbol{\sigma}_\Gamma \mathbf{n} \, ds = \int_\gamma \text{div}_\Gamma \boldsymbol{\sigma}_\Gamma \, ds$$

Interfacial surface force: $\text{div}_\Gamma \boldsymbol{\sigma}_\Gamma$

Coupling interface-fluid

$$V_\Gamma = \mathbf{u} \cdot \mathbf{n} \quad (\text{immiscibility})$$

$$[\mathbf{u}]_\Gamma = 0 \quad (\text{viscosity})$$

$$[\boldsymbol{\sigma}\mathbf{n}]_\Gamma = \text{div}_\Gamma \boldsymbol{\sigma}_\Gamma \quad (\text{momentum conservation})$$

Classical models

Interfacial effects

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Classical models

Clean interface: only surface tension

$$\boldsymbol{\sigma}_\Gamma = \tau \mathbf{P} \quad \text{with } \mathbf{P} := \mathbf{I} - \mathbf{nn}^T \quad (\text{projection on } \Gamma)$$

then

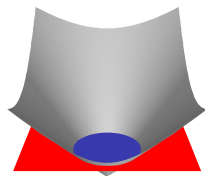
$$\operatorname{div}_\Gamma \boldsymbol{\sigma}_\Gamma = -\tau \kappa \mathbf{n} + \nabla_\Gamma \tau \quad (\kappa: \text{curvature of } \Gamma)$$

“Modeling” of the interface: Level set approach

$\Gamma(t)$ = zero-level of a scalar function:

the level set function $\varphi(x, t)$

$$\varphi(x, t) = \begin{cases} < 0 & \text{for } x \text{ in phase } \Omega_1 \\ > 0 & \text{for } x \text{ in phase } \Omega_2 \\ = 0 & \text{at the interface} \end{cases}$$



should be an “*approximate signed distance function*”.

$$x(t) \in \Gamma(t) \Rightarrow \varphi(x(t), t) = 0.$$

Level set equation

$$\varphi_t + \mathbf{u} \cdot \nabla \varphi = 0$$

Fluid dynamics model: Navier-Stokes + LS equation

Navier-Stokes equations coupled with level set equation

$$\begin{aligned}\rho(\varphi)\left(\mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u}\right) - \operatorname{div}\left(\mu(\varphi)\mathbf{D}(\mathbf{u})\right) + \nabla p &= \rho(\varphi)\mathbf{g} + \delta_\Gamma(\varphi)\operatorname{div}_\Gamma\boldsymbol{\sigma}_\Gamma \\ \operatorname{div}\mathbf{u} &= 0 \\ \varphi_t + \mathbf{u} \cdot \nabla\varphi &= 0\end{aligned}$$

Localized force term in **weak** formulation

$$f_\Gamma(\mathbf{v}) = \int_\Gamma (\operatorname{div}_\Gamma\boldsymbol{\sigma}_\Gamma) \cdot \mathbf{v} \, ds, \quad f_\Gamma \in H^{-1}(\Omega)$$

Accurate discretization of f_Γ is essential!

Mass transport (or heat transport)

$$\frac{\partial c}{\partial t} + \mathbf{u} \cdot \nabla c = D(\varphi) \Delta c,$$

$$[D(\varphi) \nabla c \cdot \mathbf{n}] = 0 \quad \text{at the interface,}$$

$$c_1 = C_H c_2 \quad \text{at the interface.}$$

D : piecewise constant. Henry condition: **discontinuity** in c .

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Transport of surfactants: PDE on $\Gamma(t)$

$$\frac{\partial S}{\partial t} + \mathbf{u} \cdot \nabla S + S \operatorname{div}_{\Gamma} \mathbf{u} = \operatorname{div}_{\Gamma} (D_{\Gamma} \nabla_{\Gamma} S) \quad \text{on } \Gamma$$

with D_{Γ} : interfacial diffusion coefficient.

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with D_{Γ} : interfacial diffusion coefficient.

No chemical reactions considered!

Numerical issues

Numerical challenges

- **Unknown** interface $\Gamma(t)$.
- Highly **nonlinear couplings** between \mathbf{u} , φ , $f_{\Gamma(\varphi)}$, c , S .
- Treatment of force balance $[\boldsymbol{\sigma}\mathbf{n}]_{\Gamma} = \text{div}_{\Gamma} \boldsymbol{\sigma}_{\Gamma}$ (**driving interfacial forces**)
- **Discontinuities** (in p , ρ , μ , c) across $\Gamma(t)$.
- PDE **on** the evolving interface.

Numerical challenges

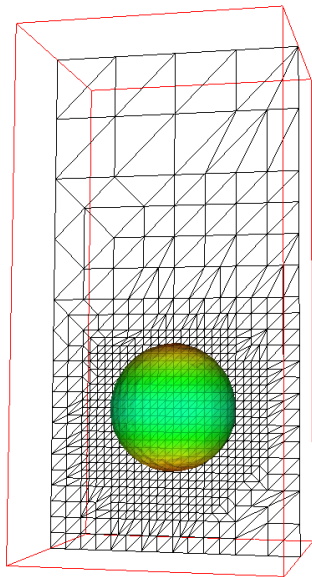
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Demand for:

- **Adaptive** space discretization, that can handle **discontinuities**.
- **Accurate** time integration, that can handle **moving discontinuities**.
- Methods for PDEs **on interfaces**.
- **Efficient and robust** iterative solvers.

Examples of numerical simulations

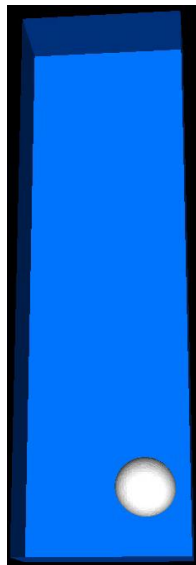
Multilevel adaptive finite elements



Rising butanol droplet: fluid dynamics

system: n-butanol/water

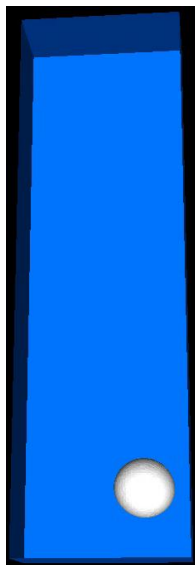
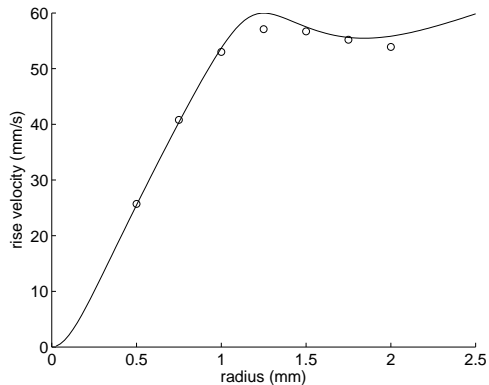
drop radius $r = 3 \text{ mm}$ \rightsquigarrow wobbling



Rising butanol droplet: fluid dynamics

system: n-butanol/water

rise velocity vs. drop radius



Rising droplet with surfactant transport

- gravity-driven butanol-droplet (diam. 4mm) in water, interfacial tension 1.63mN/m
- Velocity field determined from NS-equations.

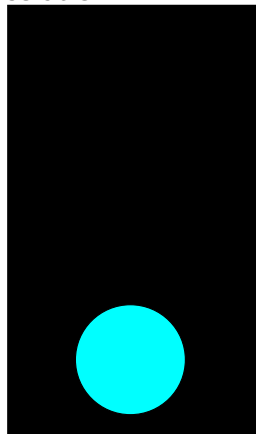
$$+ \text{surfactant eqn. } \dot{S} - D_{\Gamma} \Delta_{\Gamma} S + (\nabla_{\Gamma} \cdot \mathbf{v}) S = 0$$

- τ does **not** depend on S .

Observation

- accumulation of surfactant at the bottom.

solution



A few numerical aspects

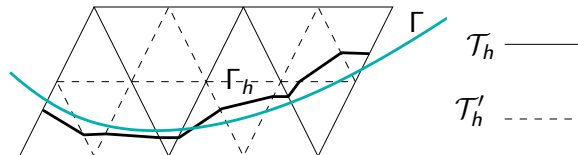
Approximation of Γ by Γ_h (hybrid level set)

Γ = zero level of φ (= level set function)

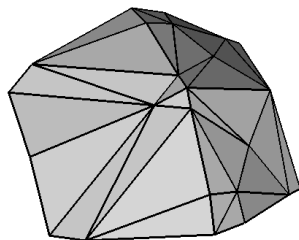
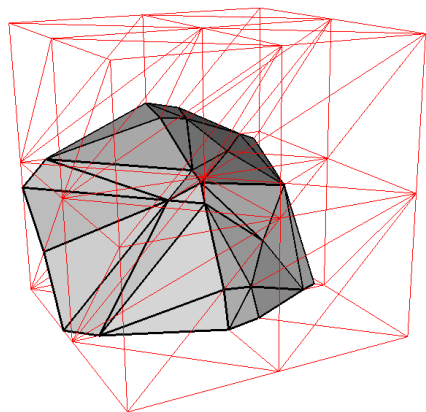
φ_h = piecewise **quadratic** FE approximation of φ .

Our strategy for Γ_h :

φ_h (piecewise P_2) \rightarrow $I(\varphi_h)$ (piecewise P_1 on refined mesh).



Approximation of Γ by Γ_h : 3D illustration



Mass transport equation

$$\frac{\partial u}{\partial t} + \mathbf{w} \cdot \nabla u - \operatorname{div}(\alpha \nabla u) = f \quad \text{in } \Omega_i, \quad i = 1, 2, \quad t \in [0, T],$$

$$[\alpha \nabla u \cdot \mathbf{n}]_{\Gamma} = 0,$$

$$[\beta u]_{\Gamma} = 0,$$

$$u(\cdot, 0) = u_0 \quad \text{in } \Omega_i, \quad i = 1, 2,$$

$$u(\cdot, t) = 0 \quad \text{on } \partial\Omega, \quad t \in [0, T].$$

With $\alpha = \alpha_i > 0$, $\beta = \beta_i > 0$.

Space-time weak formulation

Assumptions: $u_0 = 0$, $\operatorname{div} \mathbf{w} = 0$, $V_\Gamma = \mathbf{w} \cdot \mathbf{n}$ (Γ transported by \mathbf{w}).

Space-time: $Q_T := \Omega \times (0, T)$,

$$\Gamma_* := \{ (x, t) \mid x \in \Gamma(t), t \in (0, T) \}.$$

$$Q_i := \{ (x, t) \mid x \in \Omega_i(t), t \in (0, T) \}, \quad i = 1, 2.$$

Spaces (anisotropic; $Q_T = Q_1 \cup Q_2$):

$$H^{1,0}(Q_i) = \{ u \in L^2(Q_i) \mid \frac{\partial u}{\partial x_j} \in L^2(Q_i), j = 1, 2, 3 \}$$

$$H_0^{1,0}(Q_T) = \{ u \in L^2(Q_T) \mid \frac{\partial u}{\partial x_j} \in L^2(Q_T), j = 1, 2, 3, u|_{\partial\Omega} = 0 \}$$

$$V_\beta = \{ u \in L^2(Q_T) \mid u_i \in H^{1,0}(Q_i), i = 1, 2, u|_{\partial\Omega} = 0, [\beta u]_{\Gamma_*} = 0 \}$$

$$W_\beta = \{ v \in V_\beta \mid \frac{\partial v}{\partial t} \in H_0^{1,0}(Q_T)' \}.$$

Well-posed weak formulation

Determine $u \in W_\beta$ with $u(\cdot, 0) = 0$ such that

$$\frac{\partial u}{\partial t}(v) - \int_{Q_T} u \mathbf{w} \cdot \nabla v \, dx \, dt + \sum_{i=1}^2 \int_{Q_i} \alpha_i \nabla u_i \cdot \nabla v \, dx \, dt = \int_{Q_T} f v \, dx \, dt$$

for all $v \in H_0^{1,0}(Q_T)$

Note:

- Space-time formulation.
- Trial functions are discontinuous across Γ_* .
- Condition $[\beta u]_\Gamma = 0$ **essential** condition in space W_β .

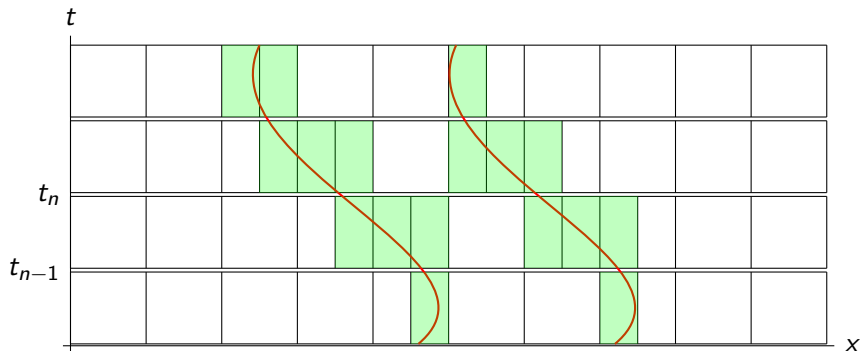
Nitsche-DG-XFEM discretization

Space-time FE.

$I_n = (t_{n-1}, t_n]$. $Q^n = \Omega \times I_n$. V_n : standard FE space on Ω .

$$W_n := \{ v : Q^n \rightarrow \mathbb{R} \mid v(x, t) = \phi_0(x) + t\phi_1(x), \quad \phi_0, \phi_1 \in V_n \} \quad (1)$$

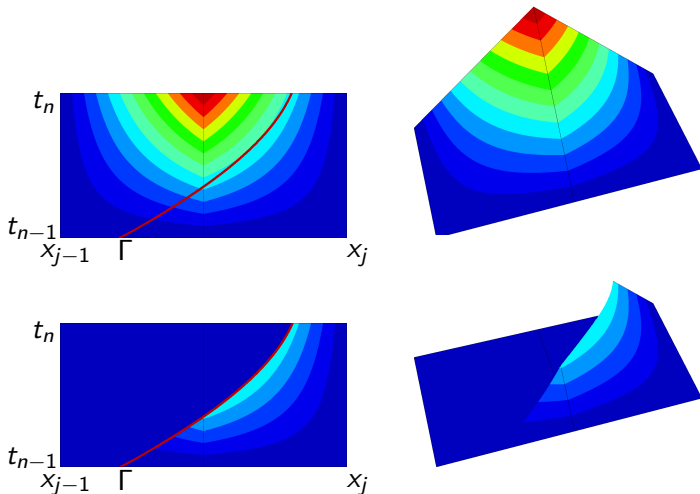
$$W := \{ v : Q \rightarrow \mathbb{R} \mid v|_{Q^n} \in W_n \} \quad (\text{space-time FE}). \quad (2)$$



Space-time XFEM.

$$Q_i^n := \cup_{t \in I_n} \Omega_i(t), \quad R_i^n : \text{restriction to } Q_i^n$$

$$W_n^\Gamma := R_1^n W_n \oplus R_2^n W_n, \quad W^{\Gamma*} := \{v : Q \rightarrow \mathbb{R} \mid v|_{Q^n} \in W_n^\Gamma\}$$



$$a^n(u, v) = \sum_{i=1}^2 \int_{Q_i^n} \left(\frac{\partial u_i}{\partial t} + \mathbf{w} \cdot \nabla u_i \right) \beta_i v_i + \alpha_i \beta_i \nabla u_i \cdot \nabla v_i \, dx \, dt$$

Discontinuous Galerkin (DG) w.r.t. time:

$$d^n(u, v) = \int_{\Omega} \beta(\cdot, t_n) [u]^{n-1} v_+^{n-1} \, dt$$

Nitsche method for Henry condition:

$$\begin{aligned} N_{\Gamma_*}^n(u, v) &= - \int_{\Gamma_*^n} \{ \alpha \nabla u \cdot \mathbf{n} \}_{\Gamma_*} [\beta v]_{\Gamma_*} \, ds - \int_{\Gamma_*^n} \{ \alpha \nabla v \cdot \mathbf{n} \}_{\Gamma_*} [\beta u]_{\Gamma_*} \, ds \\ &\quad + \lambda h_n^{-1} \int_{\Gamma_*^n} [\beta u]_{\Gamma_*} [\beta v]_{\Gamma_*} \, ds, \end{aligned}$$

with $\{ \cdot \}_{\Gamma_*}$ a suitable area weighted average. $\lambda > 0$: stabilization parameter.

$$a(u, v) = \sum_{n=1}^N a^n(u, v), \quad \text{similarly : } d(u, v), N_{\Gamma_*}(u, v).$$

Nitsche-DG-XFEM variational problem

Determine $U \in W^{\Gamma*}$ such that

$$B(U, V) = f(V) \quad \text{for all } V \in W^{\Gamma*},$$

$$B(U, V) := a(U, V) + d(U, V) + N_{\Gamma*}(U, V).$$

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Error analysis for linear FE

Theorem.

$$\|(u - U)(\cdot, t_N)\|_{L^2(\Omega)} \leq c(h^2 + \Delta t^2).$$

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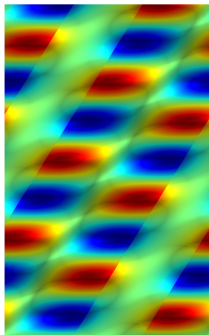
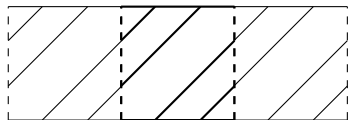
$$\|(u - U)(\cdot, t_N)\|_{L^2(\Omega)} \leq c(h^2 + \Delta t^2).$$

Remark: for standard space-time DG [V. Thomee] (no Nitsche, no XFEM):

$$\|(u - U)(\cdot, t_N)\|_{L^2(\Omega)} \leq c(h^2 + \Delta t^3)$$

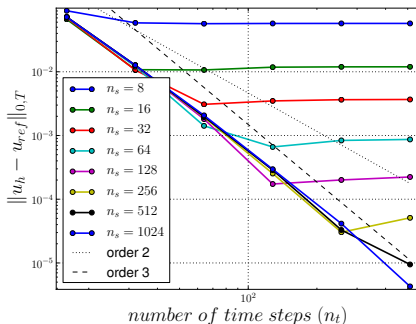
Numerical experiment

One-dimensional, linear interface velocity, periodic BC.

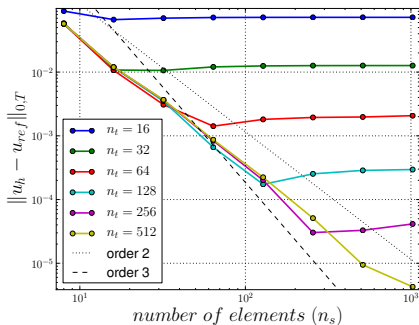


Discretization error

Error: $\|(U - u)(\cdot, t_N)\|_{L^2(\Omega)}, \Delta t = h.$



This indicates: $\|(U - u)(\cdot, t_N)\|_{L^2(\Omega)} \sim \Delta t^3$ if h sufficiently small.



This indicates: $\|(U - u)(\cdot, t_N)\|_{L^2(\Omega)} \sim h^2$ if Δt is sufficiently small.