

# Global weak solutions for an incompressible, Newtonian fluid interacting with a linearly elastic shell

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# Outline

- 1 Coupled fluid-shell system
- 2 Strategy of proof

## Koiter's energy

Koiter's energy for thin, linearly elastic shells (transverse displacements)

$$K(\eta) = \frac{1}{2} \int_{\partial\Omega \setminus \Gamma} \epsilon \langle C, \sigma(\eta) \otimes \sigma(\eta) \rangle + \frac{\epsilon^3}{3} \langle C, \theta(\eta) \otimes \theta(\eta) \rangle dA$$

Shell elasticity tensor  $C_{\alpha\beta\gamma\delta} = \frac{4\lambda\mu}{\lambda+2\mu} g_{\alpha\beta} g_{\gamma\delta} + 2\mu(g_{\alpha\gamma} g_{\beta\delta} + g_{\alpha\delta} g_{\beta\gamma})$ Linearisations  $\sigma$  and  $\theta$  at  $\eta = 0$  of the

- change of metric tensor  $\Sigma_{\alpha\beta}(\eta) = \frac{1}{2}(g_{\alpha\beta}(\eta) - g_{\alpha\beta})$
- change of curvature tensor  $\Theta_{\alpha\beta}(\eta) = h_{\alpha\beta}(\eta) - h_{\alpha\beta}$

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Taking into account inertia, the equations of motion take the form

Generalisation of the Kirchhoff-Love plate equation

$$\epsilon \rho_S \partial_t^2 \eta + \text{grad}_{L^2} K(\eta) = g \text{ in } I \times \partial\Omega \setminus \Gamma.$$

## Coupled system

## Equations of the coupled fluid-shell system

$$\rho_F(\partial_t \mathbf{u} + \nabla_{\mathbf{u}} \mathbf{u}) = \operatorname{div}(2\sigma D\mathbf{u} - \pi \operatorname{id}) + \mathbf{f} \quad \text{in } \Omega_{\eta(t)}$$

$$\operatorname{div} \mathbf{u} = 0 \quad \text{in } \Omega_{\eta(t)}$$

$$\epsilon \rho_S \partial_t^2 \eta + \operatorname{grad}_{L^2} K(\eta) = ((-2\sigma D\mathbf{u} + \pi \operatorname{id}) \boldsymbol{\nu}_t) \cdot \boldsymbol{\nu} + g \quad \text{in } \partial\Omega \setminus \Gamma$$

$$\mathbf{u} = \partial_t \eta \boldsymbol{\nu} \quad \text{on } \partial\Omega_{\eta(t)} \setminus \Gamma$$

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 \end{aligned}$$

Some important results:

**Beirão da Veiga (2004):** Local existence of strong solutions in plate case with added damping of the shell equation in 2d for small data

**Chambolle, Desjardins, Esteban, Grandmont (2005):** Global existence of weak solutions in plate case with added damping

**Grandmont (2008):** Same without damping

# Energy estimates

## Energy identity

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \int_{\Omega_{\eta(t)}} |\mathbf{u}|^2 dx + \frac{1}{2} \frac{d}{dt} \int_{\partial\Omega \setminus \Gamma} |\partial_t \eta|^2 dA + \frac{d}{dt} K(\eta) \\ = - \int_{\Omega_{\eta(t)}} |\nabla \mathbf{u}|^2 dx + \int_{\Omega_{\eta(t)}} \mathbf{f} \cdot \mathbf{u} dx + \int_{\partial\Omega \setminus \Gamma} g \partial_t \eta dA \end{aligned}$$

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By Gronwall's lemma:

$$\begin{aligned} \eta &\in W^{1,\infty}(I, L^2(\partial\Omega \setminus \Gamma)) \cap L^\infty(I, H_0^2(\partial\Omega \setminus \Gamma)) \\ \mathbf{u} &\in L^\infty(I, L^2(\Omega_{\eta(t)})) \cap L^2(I, H^1(\Omega_{\eta(t)})) \end{aligned}$$

**Note:**  $W^{1,\infty}(L^2) \cap L^\infty(H^2) \hookrightarrow C^{0,\mu}(C^{0,1-2\mu})$ ,  $0 < \mu < 1/2$



# Strategy of proof

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- Low regularity of the domain
- Hard to use difference quotients, impossible to use Aubin-Lions
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- No substitute for Bochner spaces of dual space-valued functions
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**Idea:** Give an elementary proof of Aubin-Lions theorem

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$$\int_0^t \int_{\Omega} \nabla u_n \cdot \nabla \varphi \, dx dt = \int_{\Omega} u_n(t) \varphi \, dx - \int_{\Omega} u_n(0) \varphi \, dx =: c_{\varphi,n}(t) - c_{\varphi,n}(0)$$

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$$\sup_{\|\varphi\|_{L^2(\Omega)} \leq 1} \int_I (c_{\varphi,n} - c_{\varphi}) \, dt \leq \epsilon \|u_n - u\|_{L^2(I, H_0^1(\Omega))} + c(\epsilon) \sup_{\|\varphi\|_{H_0^1(\Omega)} \leq 1} \int_I (c_{\varphi,n} - c_{\varphi}) \, dt.$$



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**We conclude:**  $u_n \rightarrow u$  in  $L^2(I, L^2(\Omega))$ .

Thank you for your attention!